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SDP bounds on quantum codes

joint works with:

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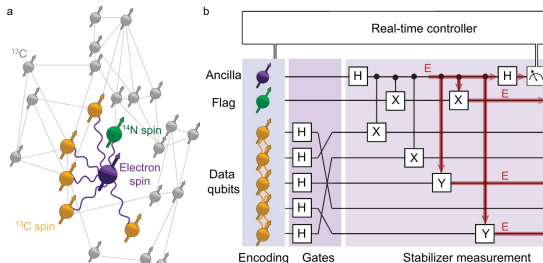
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Quantum computing



Blatt laboratory, Abobeih et al. 2022



- ▶ Information is physical (atoms, photons, electric charges ...)
 - ▶ Quantum physics is noisy / quantum information is fragile
 - ▶ Quantum information cannot be treated classically (state collapse, no-cloning, continuity of states & errors)
- ... we need a way to protect quantum states from noise

Quantum error-correcting codes

- ▶ A quantum code encodes quantum information (e.g. a qubit) redundantly. The encoded state can be recovered after being affected by noise.
- ▶ Quantum codes form the backbone of quantum computers: allow for fault-tolerant processing of quantum information.

Example

$$\alpha |0\rangle + \beta |1\rangle \mapsto \alpha |000\rangle + \beta |111\rangle$$

- ▶ circumvents no-cloning.
- ▶ discretization of errors.
(detect single flip by measurement of ZZI , ZIZ , IZZ)
- ▶ syndrome measurement collapses state onto code subspace.

Conditions for quantum error correction

- ▶ Code is a subspace $(\mathbb{C}^2)^{\otimes n}$, represented by a projector Π .
- ▶ Noise acts as a quantum channel on a state ϱ

$$\mathcal{N}(\varrho) = \sum_{E_\mu \in N} E_\mu(\varrho)E_\mu^\dagger$$

where $\sum_\mu E_\mu^\dagger E_\mu = \mathbb{1}$.

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- ▶ Knill-Laflamme conditions: $\mathcal{N}$ can be corrected on Π , iff

$$\Pi E_\mu^\dagger E_\nu \Pi = c_{\mu\nu} \Pi$$

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- ▶ Think of this as a type of **sphere packing**: equivalent to

$$\langle \phi | E_\mu E_\nu | \psi \rangle = c_{\mu\nu} \langle \phi | \psi \rangle$$

for all $E_\mu, E_\nu \in N$. and $|\phi\rangle, |\psi\rangle \in \text{ran}(\Pi)$.

Error basis

Our goal: correct all tensor-product errors of at most weight δ .

The *Pauli matrices* form a basis for complex 2×2 matrices.

$$\begin{aligned} I &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, & Y &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \\ X &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, & Z &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{aligned}$$

- Tensor-product basis

$$\mathcal{E}_n = \left\{ E_\alpha = e_{\alpha_1} \otimes \dots \otimes e_n \mid e_i \in \{I, X, Y, Z\} \right\}$$

- Weight $\text{wt}(E_\alpha)$ is the number of coordinates where E_α acts non-trivially. E.g. $\text{wt}(IXIZZ) = 3$

Fundamental question in (quantum) coding theory

A $((n, K, \delta))$ quantum code is then a projector Π on $(\mathbb{C}^2)^{\otimes n}$ of rank K , such that

$$\Pi E_a^\dagger E_b \Pi = c_{ab} \Pi$$

for all $E_a, E_b \in \mathcal{E}_n$ with $\text{wt}(E_a^\dagger E_b) < \delta$.

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- ▶ A code is *pure*, if $c_{ab} \propto \text{tr}(E_a E_b)$ for $1 < \text{wt}(E_a^\dagger E_b) < \delta$.
(maximally mixed marginals)
- ▶ A code is *self-dual* if $K = 1$ and the code is pure.

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Question

For a given block length n and distance d , what is the maximal size K for quantum codes $((n, K, d))_D$ on n quDits?

- ▶ Linear programming bounds
- ▶ Analytical bounds (q. Singleton)
- ▶ SDP bounds...?

Contributions

Result (A)

A quantum code with parameters $((n, K, \delta))_2$ exists if and only if a certain SDP hierarchy is feasible at every level.

Result (B)

The Lovász number bounds the existence of self-dual quantum codes. A symmetrization recovers the quantum Delsarte bound.

Result (C)

There a symmetry-reduced SDP of size $O(n^4)$ based on the Terwilliger algebra. Codes with parameters $((7, 1, 4))_2$, $((8, 9, 3))_2$, $((10, 5, 4))_2, \dots$ do not exist.

An attempt: Stabilizer codes

- ▶ Stabilizer group $S = \langle g_1, \dots, g_{n-k} \rangle$, $-1 \notin S$, g_i generators.
- ▶ Projector onto the code subspace

$$\Pi = \frac{1}{2^{n-k}} \sum_{i=1}^{n-k} (\mathbb{1} + g_i) = \frac{1}{2^{n-k}} \sum_{s \in S} s$$

- ▶ Simplest case: graph states $g_i = X_i \otimes_{j \in N(i)} Z_j$
- ▶ Commutative group: $[g_i, g_j] = 0$ for all g_i, g_j .

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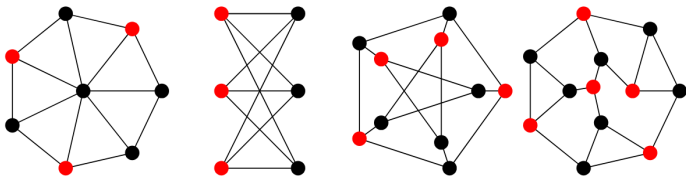
Is there a stabilizer code with $K = 1$ with distance δ (i.e. a $\delta - 1$ -uniform graph state)?

- ▶ Need maximal commuting subgroup of Pauli group.
- ▶ If $E_a E_b = -E_b E_a$, then not both can be in S .
- ▶ If $1 < \text{wt}(E_a) < \delta$, then E_a is not in S .

Independence number

Independence number (maximum size of disconnected set)

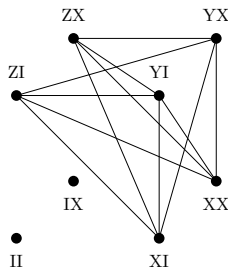
$$\alpha = \max_{W \subseteq V} |W| \quad \text{s.t.} \quad (i,j) \notin E \quad \text{for all} \quad i,j \in W$$



- ▶ It is known that $\alpha(G) \leq \vartheta(G)$
- ▶ ϑ is efficiently computable via SDP, α is not.

Lovász bound for stabilizer codes

- ▶ Index by “Pauli cube”:
 $\mathcal{E}_{n,\delta} = \{E_a \in \mathcal{E}_n \mid \text{wt}(E_a) \geq \delta\}$
- ▶ Confusability graph:
 $a \sim b$ if $E_a E_b = -E_b E_a$.
- ▶ If a self-dual stabilizer code with distance δ exists, then there is an independent set of size $\alpha = 2^n$
- ▶ As a consequence: $2^n \leq 1 + \vartheta(G)$.



Lovász bound for all quantum codes

Key idea (c.f. uncertainty relations)

- ▶ Write $\langle E_a \rangle$ for $\text{tr}(E_a \varrho)$ with $\varrho = \Pi/K$.
- ▶ Construct moment matrix $\Gamma_{ab} = \langle E_a^\dagger \rangle \langle E_b \rangle \langle E_a E_b^\dagger \rangle$ for $E_\alpha \in \mathcal{E}_n$,

$$\Gamma = \begin{matrix} & \mathbb{1} & \langle E_1 \rangle E_1^\dagger & \cdots & \langle E_N \rangle E_N^\dagger \\ \begin{matrix} \mathbb{1} \\ \langle E_1 \rangle E_1^\dagger \\ \vdots \\ \langle E_N \rangle E_N^\dagger \end{matrix} & \left[\begin{array}{cccc} 1 & \Gamma_{01} & \cdots & \Gamma_{0N} \\ & \Gamma_{11} & \cdots & \Gamma_{1N} \\ & & \ddots & \vdots \\ & & & \Gamma_{NN} \end{array} \right] & \succeq 0, \end{matrix}$$

- ▶ For two qubits, one would index with $\langle II \rangle II, \langle IX \rangle IX, \dots, \langle YZ \rangle YZ, \langle ZZ \rangle ZZ$.

Lovász bound for all quantum codes (II)

$$\Gamma = \begin{matrix} & \mathbb{1} & \langle E_1 \rangle E_1^\dagger & \dots & \langle E_N \rangle E_N^\dagger \\ \begin{matrix} \mathbb{1} \\ \langle E_1^\dagger \rangle E_1 \\ \vdots \\ \langle E_N^\dagger \rangle E_N \end{matrix} & \begin{bmatrix} 1 & \Gamma_{01} & \dots & \Gamma_{0N} \\ 1 & \Gamma_{11} & \dots & \Gamma_{1N} \\ & & \ddots & \vdots \\ & & & \Gamma_{NN} \end{bmatrix} \end{matrix} \succeq 0,$$

$$N = 4^n - 1$$

Consider $K = 1$. Then:

- ▶ $\Gamma_{00} = \langle \mathbb{1} \rangle = 1$
- ▶ $\Gamma_{ab} = \Gamma_{a0} = \langle E_a \rangle^2$
- ▶ $\sum_{a=0}^N \Gamma_{aa} = 2^n$, corresponding to $\text{tr}(\varrho^2) = 1$.

Note: If Γ is a valid moment matrix, then so is $(\Gamma + \Gamma^T)/2$.

Impose extra condition:

- ▶ $\Gamma_{ab} = 0$ if $E_a E_b = -E_b E_a$.

Lovász bound for all quantum codes (III)

This corresponds to a point in the theta body

$$\text{TH}(G) = \left\{ \text{diag}(M) \mid \begin{pmatrix} 1 & x^T \\ x & M \end{pmatrix} \succeq 0, x_a = M_{aa} \forall a, M_{ab} = 0 \text{ if } a \sim b \right\}.$$

where the *quantum confusability graph* has $\mathcal{E}_n \setminus \mathbb{1}$ as vertices and

$$\begin{aligned} a \sim b & \text{ if } \begin{cases} 0 < \text{wt}(E_a E_b) < \delta & \text{or} \\ E_a E_b = -E_b E_a \end{cases} \\ a \sim a & \text{ if } 0 < \text{wt}(E_a) < \delta \end{aligned}$$

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Lovász bound on self-dual quantum codes

If a $((n, 1, \delta))$ code exists, then

b) $\text{TH}(G)$ contains an element with $2^n = 1 + \sum_{a=1}^N M_{aa}$ and

a) $2^n \leq \vartheta(G) + 1$

Symmetry-reduction

This scales badly: Pauli cube has 4^n elements!

- ▶ Average Γ over all row and column permutations which keep triples of weights $i = \text{wt}(E_a)$, $j = \text{wt}(E_b)$, and $k = \text{wt}(E_a^\dagger E_b)$ invariant.
- ▶ The resulting matrix

$$\tilde{\Gamma} = \sum_{\pi \in \text{Aut}_0} \pi \Gamma \pi^{-1}$$

can be block-diagonalized with the Terwilliger algebra.

- ▶ This results in an SDP of size $O(n^4)$.

Gijswijt, Schrijver, Tanaka, J. Comb. Theory, A 113, 8, 2006, 1719-1731

→ Efficiently computable Lovasz bounds!

Complete SDP hierarchy for code existence

1. Formulate the Knill-Laflamme conditions $\Pi E_a E_b \Pi = c_{ab} \Pi$ as

$$K \sum_{\substack{E \in \mathcal{E}_n \\ \text{wt}(E)=j}} \text{tr}(E \varrho E^\dagger \varrho) = \sum_{\substack{E \in \mathcal{E}_n \\ \text{wt}(E)=j}} \text{tr}(E \varrho) \text{tr}(E^\dagger \varrho)$$

for $j < \delta$. (In short: $KB_j = A_j$ for $j < \delta$)

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for $j < \delta$. (In short: $KB_j = A_j$ for $j < \delta$)

2. State polynomial optimization: Consider non-commutative letters $\{x_i\}$. Form words $w = x_{j_1} \dots x_{j_k}$. Associate expectations $\langle w \rangle$ behaving as $v \langle w \rangle = \langle w \rangle v$ and $\langle v \langle w \rangle \rangle = \langle v \rangle \langle w \rangle$. State monomials have the form $w_{i_1} \langle w_{i_2} \rangle \dots \langle w_{i_m} \rangle$. Use Positivstellensatz for positive state polynomials / corresponding moment hierarchy.
(Note: Γ from above is an intermediate level!) Klep et al. 2023
- This recovers RHS in above condition. LHS...?

Complete SDP hierarchy (II)

3. Use the quantum MacWilliams identity

$$B(x, y) = A\left(\frac{x + 3y}{2}, \frac{x - y}{2}\right),$$

where $A(x, y) = \sum_{j=0}^n A_j(\Pi) x^{n-j} y^j$ and likewise for $B(x, y)$.

4. Hierarchy is dimension-free: restrict to qubits by characterization of quasi-Clifford algebras with generator relations $\alpha_i \alpha_j = (-1)^{\chi_{ij}} \alpha_j \alpha_i$, $\chi_{ij} \in \{0, 1\}$ and $\alpha_i^2 = 1$.

Gastineau-Hills 1982

5. Impose that $\varrho = \Pi/K$: use swap-like constraints,

$$\mathrm{tr}(\varrho^m) = \mathrm{tr}((1, 2, \dots, m) \varrho^{\otimes m}) = \frac{1}{K^{m-1}}$$

expanded in Pauli matrices.

Applications

- ▶ Averaging the Lovász bound over distance-preserving automorphism leads to the quantum Delsarte bound,

$$\begin{aligned} \eta = \max \quad & \sum_{j=0}^n A_j, \\ \text{subject to} \quad & A_0 = 1, \quad A_j \geq 0 \text{ with equality for } 1 < j < \delta, \\ & \sum_{i=0}^n K_j(i) A_i \geq 0 \quad \text{for } j = 0, \dots, n. \end{aligned}$$

If $\eta < 2^n$, then code does not exist.

- ▶ Hierarchy with $O(n^4)$ scaling: Average over distance and zero-preserving automorphisms. Symmetry-reduce using the 4-ary Terwilliger algebra.

Gijswijt, Schrijver, Tanaka 2006

Table

$n \setminus \delta$	2	3	4	5	6	7	8
6	16	2^α	1	0	0	0	0
7	24 - 26	2 - 3	0	0	0	0	0
8	64	8 (9)	1	0	0	0	0
9	100 - 112	12 - 13	1	0	0	0	0
10	256	24	4 (5)	0	0	0	0
11	416 - 460	32 - 42 $^\alpha$ (53)	4 - 7	2	0	0	0
12	1024	64 - 89	16 - 20	2	1	0	0
13	1586 - 1877	128 - 204	20 - 40	2 - 3	0 (1)	0	0
14	4096	256 - 295 $^\alpha$ (324)	64 - 102	4 - 9 $^\alpha$ (10)	1	0	0
15	6476 - 7606	512 - 580	64 - 138 $^\alpha$ (150)	8 - 18	1	0	0
16	16384	1024 - 1170 (1260)	128 - 256	32 - 42	4 - 6	0	0
17	26333 - 30720	2048 - 2145	512	32 - 59 $^\alpha$ (71)	4 - 8	2	0
18	65536	2048 - 4096	512 - 921 (986)	64 - 113	16 - 22	2	1
19	106762 - 123790	4096 - 7420 $^\alpha$ (8426)	1024 - 1804	128 - 249 (276)	32 - 47	2 - 3	0 (1)

Figure: Lower - upper (old bound). α : code must be impure.

For pure codes the bounds strenghten to:

$$((6, 1, 3))_2, ((11, 41, 3))_2, ((14, 290, 3))_2, ((14, 8, 5))_2, \\ ((15, 135, 4))_2, ((17, 57, 5))_2, ((19, 7314, 3))_2$$

Contributions

- ▶ Complete SDP hierarchies for bounds on quantum codes.
- ▶ Quantum analogues of the Lovász and Delsarte bounds.
- ▶ Numerically practical relaxations.
- ▶ Flexibility in applications, formally dimension-free: extensions to qudit codes & more general confusability graphs possible.

G. Munné, A. Nemec, FH, *SDP bounds on quantum codes*,
arXiv:2408.10323

G. Munné, FH, *SDP bounds on quantum codes: rational certificates*,
arXiv:2603.19901

